

Corrigendum: incompatible deformation field and Riemann curvature tensor*

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Abstract The “Corollary 1” formulation in SUN, B. H. Incompatible deformation field and Riemann curvature tensor. *Applied Mathematics and Mechanics (English Edition)*, 38(3), 311–332 (2017) is corrected. It can be stated as follows: The symmetric part of the deformation gradient has no contribution to the trace of the displacement density tensor.

Key words Riemann curvature tensor, deformation gradient, displacement density tensor

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Nomenclature

$\mathbf{X}, \mathbf{Y}$,	position vectors in the undeformed state;	∇ ,	gradient of ∇ ;
$\mathbf{x}, \mathbf{y}$,	position vectors in the deformed state;	$\nabla_{\mathbf{X}}$,	gradient of ∇ respect to $\mathbf{X}$;
$\mathbf{G}_A$,	base vector in the undeformed state;	$\nabla_{\mathbf{x}}$,	gradient of ∇ respect to $\mathbf{x}$;
$\mathbf{g}_k$,	base vector in the deformed state;	ε ,	permutation tensor;
$\mathbf{u}$,	displacement vector;	Ω ,	antisymmetric part of the displacement gradient;
$\mathbf{F}$,	deformation gradient;	$\otimes$,	tensor product;
$\mathbf{R}$,	Riemann curvature tensor in the undeformed state;	$\mathbf{I}$,	Kronecker delta, $\mathbf{I} = \delta_{ij} \mathbf{e}_i \otimes \mathbf{e}_j$;
$\Delta \mathbf{u}$,	displacement vector variation;	$\mathbf{T}$,	displacement flux density tensor.

The “Corollary 1” formulation in Ref. [1] is incorrectly presented. The flawless should be maintained, and the misrepresentation must be corrected.

1 Corollary 1 and its proof in Ref. [1]

In Ref. [1], Corollary 1 is incorrectly proposed. Making the letter self-contained, the corollary and its proof are rewritten as follows:

Corollary 1 *The symmetric part of the deformation gradient $\mathbf{F}$ has no contribution to the displacement change $\Delta \mathbf{u}$ and the displacement density tensor $\mathbf{T}$.*

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Proof Let $\mathbf{S} = \frac{1}{2}(\mathbf{F} + \mathbf{F}^T)$ and $\mathbf{\Omega} = \frac{1}{2}(\mathbf{F} - \mathbf{F}^T)$ be the symmetric part and anti-symmetric part of the deformation gradient $\mathbf{F}$, respectively. As we know, any tensor can be decomposed into a symmetric part and an antisymmetric part, i.e., $\mathbf{F} = \mathbf{S} + \mathbf{\Omega}$. Since the curl of the symmetric tensor vanishes, $\text{curl} \mathbf{S} = \mathbf{S} \times \nabla_{\mathbf{X}} = \mathbf{0}$, and $\text{curl} \mathbf{F} = \text{curl}(\mathbf{S} + \mathbf{\Omega}) = \text{curl} \mathbf{\Omega} = \mathbf{\Omega} \times \nabla_{\mathbf{X}}$. Then, we have

$$\text{curl} \mathbf{\Omega} = \text{curl} \left(\frac{1}{2}(\mathbf{F} - \mathbf{F}^T) \right) = \frac{1}{2} \boldsymbol{\varepsilon} : \mathbf{R}(\mathbf{G}_A, \mathbf{G}_B) \mathbf{u} = \frac{1}{2} \mathbf{u} \cdot \mathbf{R} : \boldsymbol{\varepsilon}. \quad (1)$$

$\text{curl} \mathbf{S} = \mathbf{0}$ indicates that the symmetric part of the deformation gradient $\mathbf{F}$ does not have any contribution towards the comparability conditions. In other words, the symmetric deformations are always compatible, and the incompatible deformation will make the symmetric deformation breaks down.

Since $\text{curl} \mathbf{\Omega} = \mathbf{\Omega} \times \nabla_{\mathbf{X}} = \mathbf{\Omega} \nabla_{\mathbf{X}} : \boldsymbol{\varepsilon}$ and thus $\mathbf{\Omega} \nabla_{\mathbf{X}} : \boldsymbol{\varepsilon} = \mathbf{R}(\mathbf{G}_A, \mathbf{G}_B) \mathbf{u} : \boldsymbol{\varepsilon}$, we have

$$\mathbf{\Omega} \nabla_{\mathbf{X}} = \frac{1}{2} \mathbf{R}(\mathbf{G}_A, \mathbf{G}_B) \mathbf{u} = \frac{1}{2} \mathbf{u} \cdot \mathbf{R}. \quad (2)$$

It must be pointed out that the statement of the above corollary and its proof are wrong and should be corrected.

2 Correction of Corollary 1 and its proof

Before presenting the corrected version of the above corollary, let us first prove a Lemma as follows:

Lemma 1 *The trace of curl of a symmetric tensor is zero.*

Proof Let $\mathbf{A} = A_{ij} \mathbf{e}_i \otimes \mathbf{e}_j$ be a symmetric tensor, i.e., $A_{ij} = A_{ji}$. The curl of the tensor $\mathbf{A}$ is defined as follows:

$$\text{curl} \mathbf{A} = \mathbf{A} \times \nabla = A_{ij} \mathbf{e}_i \otimes \mathbf{e}_j \times \nabla_k \mathbf{e}_k = \nabla_k (A_{ij}) \mathbf{e}_i \otimes \mathbf{e}_j \times \mathbf{e}_k = \nabla_k A_{ij} e_{jkm} \mathbf{e}_i \otimes \mathbf{e}_m,$$

where e_{jkm} is the permutation symbol. Therefore, we can define the trace of the curl of tensor $\mathbf{A}$ by $\text{tr}(\text{curl} \mathbf{A}) = \mathbf{I} : \text{curl} \mathbf{A} = \mathbf{I} : (\mathbf{A} \times \nabla)$, where $\mathbf{I} = \delta_{ij} \mathbf{e}_i \otimes \mathbf{e}_j = \mathbf{e}_i \otimes \mathbf{e}_i$ is the Kronecker delta. Therefore,

$$\begin{aligned} \text{tr}(\text{curl} \mathbf{A}) &= \mathbf{I} : (\nabla_k A_{ij} e_{jkm} \mathbf{e}_i \otimes \mathbf{e}_m) \\ &= \delta_{\mu\nu} \mathbf{e}_\mu \otimes \mathbf{e}_\nu : \nabla_k A_{ij} e_{jkm} \mathbf{e}_i \otimes \mathbf{e}_m \\ &= \delta_{\mu\nu} \nabla_k A_{ij} e_{jkm} (\mathbf{e}_\mu \cdot \mathbf{e}_i) (\mathbf{e}_\nu \cdot \mathbf{e}_m) \\ &= \delta_{\mu\nu} \nabla_k A_{ij} e_{jkm} \delta_{\mu i} \delta_{\nu m} \\ &= \delta_{im} \nabla_k A_{ij} e_{jkm} \\ &= \nabla_k A_{ij} e_{jki}. \end{aligned}$$

Since $A_{ij} = A_{ji}$ and $e_{jki} = -e_{ikj}$, we have $\text{tr}(\text{curl} \mathbf{A}) = \nabla_k A_{ij} e_{jki} = 0$.

Corollary 2 *The symmetric part of the deformation gradient $\mathbf{F}$ has no contribution to the trace of displacement density tensor $\mathbf{T}$.*

Proof Let $\mathbf{S} = \frac{1}{2}(\mathbf{F} + \mathbf{F}^T)$ and $\mathbf{\Omega} = \frac{1}{2}(\mathbf{F} - \mathbf{F}^T)$ be the symmetric part and the anti-symmetric part of the deformation gradient $\mathbf{F}$, respectively. As we know, any tensor can be decomposed into a symmetric part and an antisymmetric part, i.e., $\mathbf{F} = \mathbf{S} + \mathbf{\Omega}$.

Since the trace of the curl of the symmetric tensor vanishes, i.e., $\text{tr}(\text{curl} \mathbf{S}) = \text{tr}(\mathbf{S} \times \nabla_{\mathbf{X}}) = 0$, we have $\text{tr}(\text{curl} \mathbf{F}) = \text{tr}(\text{curl}(\mathbf{S} + \mathbf{\Omega})) = \text{tr}(\text{curl} \mathbf{S}) + \text{tr}(\text{curl} \mathbf{\Omega}) = \text{tr}(\mathbf{\Omega} \times \nabla_{\mathbf{X}})$. Then, we have the trace of the displacement density tensor $\mathbf{T}$ as follows:

$$\text{tr} \mathbf{T} = -\text{tr}(\mathbf{F} \times \nabla_{\mathbf{X}}) = -\text{tr}(\mathbf{\Omega} \times \nabla_{\mathbf{X}}) = -\frac{1}{2} \text{tr}((\mathbf{F} - \mathbf{F}^T) \times \nabla_{\mathbf{X}}) = -\frac{1}{2} \text{tr}(\mathbf{u} \cdot \mathbf{R} : \boldsymbol{\varepsilon}). \quad (3)$$

References

- [1] SUN, B. H. Incompatible deformation field and Riemann curvature tensor. *Applied Mathematics and Mechanics (English Edition)*, **38**(3), 311–332 (2017) <https://doi.org/10.1007/s10483-017-2176-8>